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Non-dimensional physics of pulsatile cardiovascular networks and energy efficiency

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1. Introduction

The human circulation system, with its four-chambered heart, has been rigorously studied; however, it is not the only possible working configuration present in Nature. There are a great variety of circulation systems that exhibit peculiar haemodynamics and diverse energetic performance, because of the arbitrary numbers and orientation of ventricles [1,2], auto-regulated shunts [3–5] and exotic cardiovascular valve designs [6–9] that address oxygen and nutrition transport [10–12]. For improved cardiac filling, even at ultra-low venous pressures, special mechanisms are employed, such as booster pumps [13] in the hagfish heart and stiff pericardium translating the ventricular
contraction to venous suction in the auricles of bivalve molluscs [14] and insects [15]. Both of these examples have ‘open’ circulation systems.

Despite numerous biological reports on the circulation systems of various species, major variations in their cardiovascular parameters, such as organism size, network arrangement, flow pulsatility, vascular material properties and cardiac output (CO) level, significantly challenge an elegant comparative analysis across the phylogenetic spectrum. This poses a barrier to the translation of ideas from Nature to technology, and hinders our fundamental and clinical comprehension of cardiovascular systems. Therefore, in this paper, a general unifying engineering framework is developed to compare haemodynamics, functional component designs (i.e. valves, shunts and junctions) and the energy cascades encountered in these alternative circulation system networks.

A formal hydrodynamic non-dimensionalization [16,17] of cardiovascular circuits has yet to be performed. Thus, starting with a systematic application of similitude principles, we proposed alternative non-dimensionalization schemes that address the intra- and inter-species variations of cardiovascular parameters. The present approach is based on a recent methodology, developed by our group, which covers the non-dimensional analysis of steady energy dissipation (i.e. power scaling) in any component (vessels and ventricles) of the circulatory system, and formulates the full energy budget for venous and arterial circulations [18,19]. This approach allowed disease-specific subject-to-subject comparisons and disease-to-disease evaluations by quantifying the haemodynamic severity of one vascular disease type versus another, or at different time points for the same disease [20]. In this paper, for the first time, we have formulated additional non-dimensional indices that govern vascular compliance, heart rate, wave reflections and the cardiac cycle pulsatility. Furthermore, this complete non-dimensional parameter set is integrated in the reduced-order lumped parameter network models (LPMs) of general circulation systems, making our approach independent of the cardiovascular system network.

We hypothesize that establishing the similitude conditions of the pulsatile cardiovascular flow networks will advance our understanding of cardiovascular disease states [21–23], and will influence the haemodynamic design of blood-wetted devices. For example, during cardiopulmonary bypass (CPB), the haemodynamic energy delivered to the peripheral organs can be modulated by manipulating the pulsatile flow waveform while delivering the same net perfusion flow rate through the aortic CPB cannula [24]. However, appropriate energy efficiency indices, which incorporate patient size and circulation network parameters, are not available to improve such device systems. The need for similar indices that can quantify and compare the pulsatile energetic load in cardiovascular disease [21], in the exercise performance of a patient with a single ventricle [25], in CPB [26,27], in heart valves [28,29] as well as in total cavopulmonary connection surgery [30,31], prompted the present detailed investigation. Likewise, to address clinical needs, various pulsatile flow indices have been proposed. For example, Shepard et al. [32] defined the energy equivalent pressure (EEP = total haemodynamic work/net flow rate) to quantify the energetic cost of pulsatile haemodynamic flow that differed from the steady load. Even though EEP is useful to assess the relative cost of pulsatile flow, it does not explain the physical factors that determine the pulsatile load, nor the relation of energetic cost to patient size. As the haemodynamic energy dissipation is a body size-dependent quantity, and varies significantly from patient to patient, scaling and normalization of energetic dissipation is essential for comparative clinical analysis [33].

In summary, through application of the present framework, we provide physics-based scaling indices and conduct characterization studies of the circulatory function and its associated energetic cost based on the proposed dimensionless numbers that govern pulsatile haemodynamics.

2. Methodology

2.1. Steady and pulsatile energy dissipation

The mean pulsatile energy dissipation rate ($\dot{e}_p$) is the difference between the mean total energy dissipation rate ($\dot{e}_T$) and the energy dissipation rate of the steady component of the flow ($\dot{e}_s$),

$$\dot{e}_p = \dot{e}_T - \dot{e}_s = \frac{1}{T} \int p(t) q(t) dt - \frac{1}{T} \int p(t) ~dq(t) ~dt,$$

where $p(t)$ and $q(t)$ are instantaneous pressure and flow rate at the junction of the aorta and ventricle, respectively, and $T$ is the duration of the cardiac cycle. The vascular structural properties and network configuration determine the pressure-flow waveforms in each vascular segment through pulse wave propagation and damping, which consequently determine the energetic load.

2.2. Similitude of pulsatile flow in compliant vessels

A theoretical analysis of haemodynamic power loss under non-pulsatile flow conditions has recently been provided by Dasi et al. [18,19]. Expanding this formulation to pulsatile flow regimes, we start by describing a vascular ‘compartment’ as a physical model of an isolated segment of the flow system in which the flow is governed by the physical properties and geometry of the flow domain, in addition to the physical conditions imposed at its boundaries (figure 1). In §2.2.1, the non-dimensional formulation of the flow dynamics in a distributed model of the vascular compartment is presented (figure 1a). In addition, an equivalent non-dimensional number set, which is motivated by the circulatory lumped parameter networks, is developed in §2.2.2 (figure 1b). In §2.2.3, we define a power

![Figure 1](https://example.com/figure1.png)

**Figure 1.** (a) A general three-dimensional one inlet–one outlet pulsatile vascular component is described by its physical parameters: fluidic density $\rho$, viscosity $\mu$, wall elastic modulus $E$, characteristic length $l$ and the flow rate $Q(t)$. (b) The same vascular segment described in the lumped element domain. Elastic, inertial and viscous properties are represented by compliance $C$, inertance $L$ and resistance $R$. The characteristic length, $l$, is related to the body surface area of the species, as described in the text.
loss index (PLI), which weighs the power loss in a vascular segment against a subject-specific power scale.

2.2.1. Non-dimensionalization of distributed vascular component models

The physical parameters defining a vascular compartment are: blood density \( \rho \), dynamic viscosity \( \mu \), vascular elastic modulus \( E \), characteristic length \( l \) and its geometry, which is represented by a dimensionless form vector \( S \) containing ratios of lengths, as shown in figure 1a. At the inlet boundary, a pulsatile flow \( Q(t) \) is imposed, which is characterized by its mean \( Q \) fundamental frequency (i.e. heart rate (HR)—in units of frequency) and waveform shape (spectral components of the waveform are included in a dimensionless form vector \( Q_{\text{e}} \) electronic supplementary material, appendix A). Outlet boundary conditions (BCs) can range from simple constant pressure outlets to more complex windkessel (WK) outlets, one-dimensional distributed distal flow models or three-dimensional models. In the present formulation, all parameters related to outlet BCs are contained in a dimensionless ‘\( B \)’ vector (e.g. \( \beta = P_{\text{in}}/(\rho Q^2/l^4) \) for a constant pressure outlet).

The functional dependence of local pulsatile haemodynamic power loss \( \tilde{\varepsilon}_p \) on these physical independent variables can be re-posed as a relationship between seven non-dimensional parameters related to outlet BCs are contained in a dimensionless \( \gamma \) appearing in the formulation to relate the local geometrical information (e.g. the constriction in a stenotic artery) to the changes in \( R, l \) and \( C \), if this relation is available.

2.2.2. Non-dimensionalization of lumped vascular component models

While the previous non-dimensional number set governs the pulsatile flow physics for a compartment, completely, an alternative similarity set would be more amenable for circulation network analysis. Particularly for complex vascular systems consisting of many vessels and vascular components, it may not be practical to individually consider the contribution of each element to the vascular function. Instead, the lumped haemodynamic parameters, resistance \( (R) \), compliance \( (C) \) and inertance \( (L) \), incorporate the effect of material properties (viscous, elastic and inertial, respectively) and the prevailing geometry in a single parameter that represents the haemodynamic function of the specific vascular segment that is under consideration (figure 1b). Our intention here is to develop ‘non-dimensional’ circulation networks and to be able to compare them through numerical LPMs. While higher-order analysis methods of cardiovascular systems do exist, we find LPMs to be the most practical approach, leading to a system-level understanding [34].

For a lumped compartment, the functional dependence of the mean pulsatile energy dissipation on the independent lumped variables is represented by \( \tilde{\varepsilon}_p = f(R, C, L, HR, Q, Q_{\text{w}}, S, \beta) \). Following non-dimensionalization (electronic supplementary material, appendix C), the functional dependence of \( \tilde{\varepsilon}_p \) can be re-posed as a function of six dimensionless quantities:

\[
\frac{\tilde{\varepsilon}_p}{RQ^2} = f(\delta, \psi, Q_{\text{w}}, S, \beta),
\]

where

\[
\delta = \frac{1}{HR \times RC}, \quad \psi = HR\sqrt{L/C}.
\]

Equation (2.3a) gives the ratio of pulsatile power loss to the steady power loss \( (= RQ^2) \), as a function of \( \delta \), the pulse decay number, and \( \psi \), the wave propagation number. \( \delta \) is based on the elastic and viscous components of a vascular segment, and governs the rise and decay characteristics of pressure in RC circuits, similarly to the WK model of circulation. \( \psi \) is based on the elastic and inertial components of a vascular segment, and governs the transmission of pulse waves across a given vascular segment that is characteristic of the transmission line model of the arterial system. Although the geometry is intrinsic to lumped parameters, \( S \) is kept in the formulation to relate the local geometrical information (e.g. the constriction in a stenotic artery) to the changes in \( R, L \) and \( C \), if this relation is available.

2.2.3. Pulsatile power loss index: scaling with body size and cardiac output

Scaling of power loss with the body surface area (BSA) is critical for understanding the relation of energetics to body size and metabolism, which is imperative in biological and clinical contexts. In general, human and mammalian circulation exhibit geometrical similarity, and it has been shown that \( l = \sqrt[13]{\text{BSA}} \) is an acceptable approximation for all major vessels [35,36]; thus, the characteristic length for the circulation is chosen as \( \sqrt{\text{BSA}} \). Similarly, the characteristic flow rate for the whole circulation is the CO. We define the PLI in any vascular compartment as the ratio of \( \tilde{\varepsilon}_p \) to the characteristic inertial power available in the circulation system, which is based on \( \sqrt{\text{BSA}} \) and \( \text{CO} \), by manipulating equations (2.2a) and (2.3a), respectively, giving

\[
\text{PLI} = \frac{\tilde{\varepsilon}_p}{\rho \text{CO}^2/\text{BSA}^3} = \alpha^3 \left( \text{BSA}^3/\text{CO}^4 \right) f(Re, Ca, St, Q_{\text{w}}, S, \beta),
\]

and

\[
\text{RI} = \alpha^2 R \text{RI}^2 (\delta, \psi, Q_{\text{w}}, S, \beta).
\]

where \( \alpha = Q/\text{CO} \) is the fraction of local flow to the total available flow and \( \text{RI} = R/(\rho \text{CO}^2/\text{BSA}^3) \) is the resistance index. In this report, the PLI is used as the measure of pulsatile energetic performance. In our earlier study, a ‘steady flow’ version of the PLI where \( \tilde{\varepsilon}_p \) replaces \( \tilde{\varepsilon}_p \), has been used to assess total cavo-pulmonary conduit performance [19,25]. For presentation purposes, and in order to obtain practical clinically relevant value ranges, the PLI and RI are multiplied by \( 10^{-9} \) in this paper [18].

2.3. Arterial model

To experiment with these new non-dimensional indices, we developed a numerical model of the systemic arterial circulation. Figure 2 depicts the model of systemic arterial circulation that was used in parametric simulations to determine the changes in PLI with respect to the cardiovascular non-dimensional parameters. The model consists of a one-dimensional transmission line model of the aorta that terminates with a three-element windkessel (WK3) model of the peripheral microvasculature. \( C_a \) and \( L_a \) are the aortic compliance and inertance, \( C_p \) and \( R_p \) are the peripheral compliance and resistance, and \( Z_L = \sqrt{L_a/C_a} \) is the characteristic impedance of the WK3. Arterial model parameters match the characteristics of adult
human circulation (table 1) [37]. The arterial segment is uniform; therefore, wave reflections occur only at the intersection of the arterial segment and the WK3. This is a simplification of the actual system, which contains spatially distributed wave reflections; however, the focal point of combined wave reflections can be approximated as a single site [38]. To represent a healthy physiology, the characteristic impedance of the WK was matched to the actual system, which contains spatially distributed wave reflections [42], because $Z_C$ does not show up in the non-dimensional numbers set, as it is not an independent parameter however, their investigation is outside the scope of the present study. Considering the above, the functional dependence of the cardiovascular diseases [21] and surgical reconstructions [20]; geometrical effects introduced by stenoses, aneurysms and congenital defects are significant for haemodynamic losses in many cardiovascular diseases [21] and surgical reconstructions [20]; however, their investigation is outside the scope of the present study. Considering the above, the functional dependence of the PLI on non-dimensional parameters is given in

$$PLI = R_I J(\delta, \psi).$$

As illustrated in this section, the proposed framework can be applied to any chosen cardiovascular compartment subset in a similar fashion.
Table 2. Allometric exponential coefficients of cardiovascular systems.

<table>
<thead>
<tr>
<th></th>
<th>surface-area scaling</th>
<th>quarter-power scaling</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>human [35]</td>
<td>mammals [43,47]</td>
</tr>
<tr>
<td>aortic pressure</td>
<td>0</td>
<td>0</td>
</tr>
<tr>
<td>cardiac output</td>
<td>2/3</td>
<td>3/4</td>
</tr>
<tr>
<td>heart rate</td>
<td>−1/3</td>
<td>−1/4</td>
</tr>
<tr>
<td>stroke volume</td>
<td>1</td>
<td>1</td>
</tr>
<tr>
<td>characteristic velocity</td>
<td>0</td>
<td>0</td>
</tr>
<tr>
<td>vascular dimensions</td>
<td>1/3</td>
<td>3/8</td>
</tr>
<tr>
<td>blood volume</td>
<td>1</td>
<td>1</td>
</tr>
<tr>
<td>total vascular resistance</td>
<td>−2/3</td>
<td>−3/4</td>
</tr>
<tr>
<td>total vascular compliance</td>
<td>1</td>
<td>−3/4</td>
</tr>
<tr>
<td>total arterial inertia</td>
<td>−1/3</td>
<td>−1/2</td>
</tr>
</tbody>
</table>

Note: Allometric functions based on body surface area (BSA) are converted to body volume relations by multiplying the allometric exponent for BSA by two-thirds under the observation that BSA scales with two-thirds of the body weight.

The exponential coefficients for lumped parameters are estimated from the allometric relations reported in [35,43,46].

3. Results
3.1. Validation of pulsatile similitude via allometric relations

Allometric relations are statistically obtained power-law relations of the form $Y = a b^b$ that relate the value of a physiological parameter ($Y$) to body size ($b$) with a scaling exponent $b$ and a normalization constant $a$. Allometric relations do not explain the variation of a variable for subjects with the same body sizes, but as it represents the statistical average of a large population, an allometrical trend can be interpreted as a ‘normal’ ball-park value for a physiological measure as a function of body size [33].

It is widely accepted that the relation between the size and form of efficient transportation networks dictates the allometric scaling of physiological variables with body mass (i.e. $Y \sim M^b$), as shown through analytical models [43–46] and supported by empirical studies [35,47]. Observed values of $b$ generally appear as multiples of 1/4 and 1/3, and are, respectively, attributed to the fractal and Euclidean scaling of geometry in biological organisms [48–50].

In this section, we propose that non-dimensionalized circulatory systems should yield scale-invariant characteristics, if their design is governed by a common optimality principle. For this purpose, we tested the proposed non-dimensional parameters for similitude (scale invariance) by the application of allometric relations for mammals (table 2) [35,43,45–47]. Two scaling schemes are considered in this study: quarter-power and surface-area scaling schemes, which are based on the empirically observed scaling of basal metabolic rate with body mass increased to the power of three-quarters and two-thirds; the former is commonly accepted in biological sciences, whereas the latter is accepted in the clinical field. Parameters not included in table 2 are the haemorheological properties (viscosity and density) and wall elastic modulus, which do not show a significant association with body size [48–50]. Normalized flow waveform shape, given by $Q_{oc}$, is considered to be independent of body size, as normalized ventricular contraction patterns and arterial flow waveforms are closely similar in mammals [51].

Allometric equations for similitude parameters governing local flow dynamics in a single conduit vessel component can be computed as

$$Re \sim M^{1/3}, \quad Ca \sim M^0 \quad \text{and} \quad St \sim M^0 \quad \text{(surface-area scaling)}$$

$$Re \sim M^{-1/2}, \quad Ca \sim M^{-1/6} \quad \text{and} \quad St \sim M^0 \quad \text{(quarter-power scaling)}.$$  

Although there is geometrical similarity in the large vessels of an animal (i.e. $S \sim M^0$), as expected, only the dynamic similarity is not conserved, as inertial forces become dominant over viscous forces as the subject size increases.

Scaling of a single vessel gives an incomplete picture for the scaling of circulation as a system. To explore similitude in the entire circulation network, we refer to the scaling of the arterial system as a whole. We direct our attention to the most important aspects of pulsatile flow: frequency-dependent pulse propagation, damping and reflection, which can be adequately captured by the arterial model explained in §2.3. Allometric scaling of resistance, compliance and inertia can be determined empirically, or can be derived from the models of the underlying vascular structures (table 2). The body size proportionality of biological flow indices is determined from the above-mentioned allometric relations as:

$$RI \sim M^0, \quad \delta \sim M^0 \quad \text{and} \quad \psi \sim M^0 \quad \text{(surface-area scaling)}$$

$$RI \sim M^{-1/6}, \quad \delta \sim M^{1/6} \quad \text{and} \quad \psi \sim M^0 \quad \text{(quarter-power scaling)}.$$  

The exponential coefficients for $\delta$ and $\psi$ are zero, indicating that the non-dimensional parameters governing the pulsatile system are size invariant. It should be noted that the exponential scaling of the RI in the quarter-power scaling case does not break the generality of this finding, as it is not essentially a similitude parameter, but rather adjusts the proportionality of the PLI. As such, the exponential coefficient for the PLI is expected to be less than 0 (surface-area scaling), and greater than or equal to $−1/6$ (quarter-power scaling).
Therefore, the presented results confirm the validity of system-level similitude of the mammalian circulatory system when, regardless of the empirical scaling law options, non-dimensional parameters are formed with appropriately chosen length and time scales. The size-invariance property of non-dimensional characteristics enables the direct comparisons of cardiovascular operational and performance states to be made between different subjects.

3.2. Generalized pulsatile flow characteristics of the arterial system

In the present section, the variability in energetic characteristics that may accompany deviations from normal states is inspected as a function of similitude parameters using the parametric numerical model of the arterial system, as described in the Methodology (§2.3).

3.2.1. Pulsatile power loss index as a function of propagation and decay numbers

In this section, we inspect the impact of variations in propagation and damping characteristics governed by \( \psi \) and \( \delta \) on the PLI, while isolating these effects from the mean resistive effects by fixing the value of the RI during simulations. The functional dependence of the PLI on \( \delta \) and \( \psi \), while the RI is held constant at 3.85, is displayed as a surface graph in figure 3. The ‘baseline state’ is defined as the arterial model for a healthy adult human, as described in §2.3. At this baseline state of the circulatory system, \( \delta \) and \( \psi \) are 0.7 and 0.074, respectively, and the resulting PLI is 0.6. \( \delta \) and \( \psi \) are both varied between 25% and 400% of their baseline values, and changes in the PLI are plotted on the state surface of figure 3. The PLI increases monotonically with both \( \delta \) and \( \psi \). Alternately doubling and then halving \( \delta \), while keeping the other non-dimensional parameters constant, resulted in a PLI of 1.4 (+127% change) and 0.3 (−54%), respectively. The same changes in \( \psi \) resulted in a PLI of 1.8 (+200%) and 0.3 (−50%), respectively. Simultaneously increasing or decreasing \( \delta \) and \( \psi \) amplified the change in the PLI. No local extrema were observed in the power loss function.

The effects of changes in a single parameter (HR, \( R \), \( L \) or \( C \)) are also plotted in figure 3. It was found that PLI increased with an increase in stiffness or inertance. PLI also increased with either a decrease or an increase in HR, indicating that a minimal PLI exists at an intermediate HR. An increase in \( R \) led to a decrease in PLI due to increased WK performance, even when the RI is allowed to increase. However, the total power loss significantly increased because of the steady component. The impact of changes in individual vascular properties over the PLI is further discussed in §4.1 and summarized in table 3.

One case that is worthy of investigation is the dependence of the PLI on the variability of the heart rate while the parameters depending on vascular properties (\( R \), \( L \), \( C \)) are held constant. This case is clinically interesting because the heart rate is a typical, acutely controlled system parameter, as opposed to the vascular geometry and material properties, which can be altered through long-term remodelling and growth processes. This observation leads to the constraint of acutely fixed vascular states, where only changes in heart rate are allowed. These iso-contours are displayed in figure 3, on curves defined as \( \delta \times \psi = \sqrt{L/C_0}/R = k \), where \( k \) is a constant. The PLI is observed to increase with \( k \) when \( RI \) is unchanged. Under the constraint of fixed vascular state, an optimal (\( \delta \), \( \psi \)) pair exists for each ‘\( k' \), at which point the PLI reaches a local minimum. As observed, the simulated baseline (i.e. healthy) circulatory model was very close to the optimal condition,
Table 3. Influence of vascular properties on expected power loss through linear and nonlinear mechanisms. Combination of linear effects (propagation, damping, mean resistance), estimated by the arterial model, are given as overall influence on power loss. Sign of arrows indicates increasing or decreasing values; the implication for power loss is stated as favourable/unfavourable for a given flow rate.

<table>
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<th>parameter</th>
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<th>implication for power loss</th>
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variation in vascular properties

3.2.2. Arterial pressure waveform as a function of propagation and decay numbers

Figure 5 shows the change in the normalized pulsatile pressure waveform at the inlet with respect to changes in $\delta$ and $\psi$. An increase in $\delta$ is reflected in the separation of systolic and diastolic pressures, indicating that the damping factor of the WK property is diminished (figure 5a). The slope of the diastolic pressure decay strongly correlates with the $\delta$ number. The arrival time of backwards travelling waves, originating from the arterial–peripheral boundary, increases in relation to the propagation number $\psi$, but not to $\delta$. Increasing $\psi$, while keeping $\delta$ and RI constant, led to the observation that early arterial pulse pressure increases in relation to aortic impedance ($=\sqrt{L_{a}/C_{a}}$), which relates the instantaneous pressure and flow rate of the arterial pulse wave, thereby increasing the power requirement by the ventricle (figure 5b). It is also observed that the arterial model increasingly behaves as a classical RC-WK as $\psi$ decreases [55].

The occurrence of an optimal state at an intermediate ($\delta, \psi$) pair value, when the vascular state is fixed and either $\psi$ or $\delta$ is varied, is also evident in figure 5c: at a high $\psi$ (low $\delta$), WK pressures are low but high pressures are required for generating the forward travelling pulse; on the contrary, WK losses are significant at low $\psi$ (high $\delta$).

3.3. Compartmentalization and non-dimensionalization of a general circulatory network

While the dimensionless relations shown in equations (2.2)–(2.4) are general, their physical role in a multi-component circulatory network system needs to be clarified. In our approach, any circulatory network can be divided into imaginary compartments and the connection map among these compartments can be determined. For example, figure 6 showcases the ‘compartmentalization’ scheme of sample complex circulation networks: crocodile and octopus [56–58]. Individually, each of these compartments may consist of one or more lumped elements, hierarchical, multi-scale, one-dimensional or three-dimensional distributed models. There is no restriction on the model being open or closed. In open models, BCs, and in
closed models, activation functions (e.g. ventricular time-varying elastance function), should be given as inputs to the non-dimensionalization function. The final step is the non-dimensionalization of the parameters using the conventions presented in §2.2. Alternatively, for an \( N \) compartment circulation network, there are at least \( N \) coupled differential equations to solve, after which non-dimensionalization of each equation will lead to the same compartmental non-dimensional number set (not shown here for brevity).

The total ventricular load (\( \dot{E}_V \)) can be obtained from the sum of the cycle-averaged integral of the pressure–volume loop of each ventricle,

\[
\dot{E}_V = \sum_i \frac{1}{T} \int P_i \, dV_i, \tag{3.1}
\]

where \( V \) is the ventricular volume and summation is over all ventricles, as labelled by \( i \). All the power produced by the ventricles is eventually dissipated as heat at the components of the circulatory system, which can be determined individually for each compartment using equations (2.4a,b). Application of the energy conservation principle, combined with equations (2.4a,b), determines the energy budget for

\[
\frac{e_{\text{pulse}}}{e_{\text{total}}} = \frac{X}{1 + \frac{\pi}{2} \sum_i V_i}, \tag{3.3}
\]
the entire cardiovascular system as:

$$PLI_v = \sum_{\text{all compartments}} PLI_i$$

$$= \sum_{\text{lumped}} \alpha_i^2 R_{I_i} f_i(\delta_i, \psi_i, Q_{w,i}, S_i, \beta_i)$$

$$+ \sum_{\text{distributed}} \alpha_i^3 \left( \frac{l_i}{\text{BSA}} \right)^{-4} g_i(R_{c,i}, C_{a,i}, S_t, Q_{w,i}, S_i, \beta_i),$$

(3.2)

where each compartment is labelled by subscript $i$. ‘Lumped’ and ‘distributed’ indicate that the compartments are defined by the corresponding dimensionless number sets, given by equations (2.2) and (2.3).

In equation (3.2), BCs are assumed individually for all compartments. In fact, if the network connectivity map is known, such as for the circulations in figure 6, with all compartmental information present, then it is sufficient to know the BCs at open boundaries only (there would be no open boundaries for closed networks). For connected compartments, BCs at junctions with neighboring compartments (i.e. $\beta$) and the distribution of flow (i.e. $\alpha$) will be determined by the properties of the rest of the compartmental network, to which the compartment under investigation is connected, and by the BCs imposed from the open boundaries of the largest connected subgraph. A definitive implication is that local changes in compartmental properties will have global effects. As an example, suppose a network has two compartments, such as arterial and peripheral as shown in §3.1, an increase in stiffness in the arterial compartments would affect power loss in the periphery by altering the flow–pressure waveforms imposed on the microvasculature.

As the number of compartments increases along with increasing model complexity, the quantity of non-dimensional numbers that are required to represent the entire network also increases, because of a larger number of degrees of freedom. Therefore, determining the least number of parameters that can sufficiently define the compartment models, through model reduction and importance analysis, is critical for non-dimensional network analysis [59,60].

4. Discussion

4.1. Energetic implications of non-dimensional parameters

In our earlier studies, it has not been possible to validate the steady non-dimensional energy dissipation scheme of Dasi et al. [18] against the empirical allometric relations because of the incomplete non-dimensional number set. This study completed the governing number set by providing a unified picture of the transmission and WK effects on the pulsatile energetic load, and showed that the present generalized non-dimensional set is valid for both steady and pulsatile flow regimes, providing an indirect physical explanation for the cardiovascular allometric observations through physical similitude for the first time in the published literature.

Pulse propagation and reflection characteristics in mammals have been investigated empirically using allometric relations and non-dimensional parameters and have been found to be independent of animal size; however, an explanation for this observation was not previously provided [61]. In a series of papers, Pahlevan & Charib [62,63] investigated the impact of arterial stiffness, heart rate and wave reflections on the pulsatile power loss. Their study found
optimality criteria for mammals based on matching the heart rate with the travel time of pulse waves along the aorta to minimize ventricular load [64]. They proposed an invariant wave condition number, which is equal to the product of the heart rate and the effective length of the aorta divided by the pulse wave velocity (PWV), which in turn is equal to 0.1 when the heart rate is optimized. If the effective length is close to the aortic length, as is the case for healthy subjects [65], then the invariant number from Pahlevan and Gharib corresponds directly to the $\psi$ number, as $1/\sqrt{LC}$ is the natural frequency of the aorta and is equal to PWV divided by the aortic length. To account for cases in which the effective length of the aorta differs from its anatomical length, the ratio of the effective length to the anatomical length can be incorporated into $S$; this allows Pahlevan and Gharib’s wave condition number to be obtained by combining $S$ and $\psi$. Unlike these previous studies, which neglected the damping characteristics of the peripheral circulation, our analysis demonstrated that the $\psi$ number alone does not guarantee optimal cardiac performance, unless the vascular state is fixed. We found that, for each vascular state, the optimal $\psi$ value changes, as well as the corresponding $\delta$ value, although the former varies less than the latter.

It is well known that an increase in vascular stiffness impairs pulse buffering by the WK effect and increases the pulse pressure faced by the ventricle [66]. Westerhof & Elzinga [67] reported their observations on the invariance of the product of the heart rate and the arterial pulse decay time of WK3 based on empirical measurements in mammals; their data suggest that $0.2 < \delta < 0.5$. Our study demonstrates that the arterial system operates more efficiently when $\delta < 1$, which suggests that the mammalian peripheral vascular system is adjusted to minimize the pulsatile workload of the ventricle.

Physically, a constant $\psi$ indicates that the spatial and temporal distribution of forward and reflected waves is similar across species, and the constancy of $\delta$ implies that the pressure decay during a cardiac cycle is identical. In sum, it is implicitly assumed that the pressure wave faced by the heart during the ejection period is similar for all mammals. This is generally true for most mammals; one exception that is frequently mentioned is the kangaroo, which has an unusual pressure waveform [68]. An inspection of pressure and flow tracings from non-mammalian species showed waveform similarities between tuna [69], turtle [70], alligator [56] and mammals. Through a WK2 parametric estimation, we calculated a $\delta = 0.34 \pm 0.12$ across vertebrate species, where $\delta \sim 0.45$ for active species such as mammals and tuna and $\delta \sim 0.25$ for the less active turtle and alligator. Past experiments on lower vertebrates (toad, lizard and snake) have shown that pulse propagation plays a minor role compared with its role in active animals, as the pulse transit time-to-cardiac cycle duration ratio is significantly lower (5–10 times lower $\psi$ than in mammals, according to [71]) and their arterial system acts mainly as a WK [71]. We suggest that the observed invariance of $\delta$ could be valid for lower vertebrates and mammals.

**4.2. Implications for achieving optimized states beyond the baseline**

From figure 3, it is observed that an additional reduction of pulsatile power loss from the baseline state is possible by decreasing $L$ or increasing $C$, but these states may not be achieved because of mechanical or physiological design...
constraints, such as the limitation on allowed vascular volume or nonlinear dynamic effects. Therefore, in table 3, we summarize the impact of different vascular properties on the power-loss mechanisms governed by linear flow effects through the main findings of this paper, and discuss nonlinear effects that are not included in the linear model, such as steep wavefronts and turbulence. Considering that $\rho$ and $\mu$ are relatively constant, increasing aortic luminal area ($A_0$) or decreasing vascular elastic modulus represent the available options to decrease power loss by obtaining a smaller $R$, $L$, and a larger $C$, as traced in figure 3. However, increasing $A_0$ requires a larger blood volume; while, on the other hand, decreasing $E$ might lead to a low PWV, resulting in the occurrence of nonlinear steep wavefronts, if the flow velocity approaches PWV. Therefore, these two factors may impose a lower limit to the reduction of pulsatile energetic load. Strain-stiffening elasticity is another nonlinear effect, but it only contributes to a higher power loss when blood pressure is high, which necessitates an already high power loss. Thus, it is not expected to constitute a limiting factor in the low-power design.

4.3. Relation between lumped and distributed non-dimensional parameters

Equations (2.2) and (2.3) are applicable in both steady and pulsatile flow regimes. Under steady flow, $\delta$ and $\psi$ as well as $Q_m$ vanish, and the PLI becomes a function of $RI$ and $S$ alone (pulsatile power loss is replaced by steady power loss). The framework proposed in Dasi et al. [18] and ours, proposed here, can be linked using the Darcy–Weisbach equation for pressure loss, which is a function of the Darcy friction factor $f_D$, pipe length $L$ and diameter $d$, mean flow velocity $u$ and fluid density $\rho$. $\Delta P = f_D \times \frac{L}{d} \times \frac{u^2}{2}$. As resistance is the ratio of mean pressure drop to the mean flow rate, its functional dependence is identical to the Darcy–Weisbach equation for pressure loss: $\Delta P/Q = R = f (f_D, L, D, u, \rho)$. Non-dimensionalization of this expression returns $RI$ and its functional dependence on $f_D$ and $L/D$ as $R/(\rho Q/D^4) = f (f_D, L/D)$. Considering pipe flow, $f_D$ can be determined empirically from the Moody chart, if the Reynolds number and the surface roughness are known. Thus, the definition of PLI as a function of $RI$ is identical to its definition based on $Re$: $PLI = f (RI, S) = f (Re, S)$.

As explained in §4.1, $\psi$ is equivalent to the ratio of arterial length to the distance travelled by a pulse wave in the duration of a cardiac cycle

$$HR\sqrt{LC} = \frac{HR}{PWV}.$$  

PWV can be estimated with the Moens–Korteweg formula: $PWV = \sqrt{\frac{Eh}{pd}}$, where $h$ and $d$ are wall thickness and internal diameter [39], respectively. Thus, the $\psi$ number is identical to

$$\frac{HR}{\sqrt{Eh/pd}} = \sqrt{\frac{St}{Ca}} \sqrt{h/d},$$

with characteristic length chosen as the longitudinal vascular length and $h/d$ is the shape number. The $\delta$ number governs the frequency-dependent interaction of viscous fluid and elastic wall forces. Combining $Re$, $Ca$, $St$ and $S$ numbers gives an identical relation:

$$\delta = \left(\frac{E}{HR\mu}\right)^{1/2} = \left(\frac{Re}{Ca}\right)^{1/2} \left(\frac{St}{S}\right)^{1/2}.$$  

This emphasizes the compatibility of the proposed lumped parameter description of power loss to its distributed parameter description that was derived in §2.2.

4.4. Inertial and viscous scaling

From equations (2.4a,b), it is observed that mean power loss should be proportional to the cube of the flow rate in order for PLI to remain insensitive to variations in the mean flow rate. This relation is valid for cardiovascular segments where energy dissipation is dominated by inertial losses under high-$Re$ conditions. However, under laminar flow regimes ($Re < 2000–3000$), viscous forces dominate energy dissipation and the power-loss dependency on flow rate becomes a quadratic function of the flow rate ($c \approx Q^2$). It is important to make this distinction, as most power loss in the cardiovascular system occurs at the level of the microcirculation, where viscous forces are dominant. Equations (2.4a,b) can be adjusted to a suitable form for flow regimes that are dominated by viscous forces by simply multiplying the left-hand side by the Reynolds number as:

$$\frac{ep}{\mu (Q^2/BSA^{3/2})} = \frac{RL_m}{f(\delta, \psi, Q_m, S, \beta)} = \frac{Reg(Re, Ca, St, Q_m, S, \beta)}{\mu (Q^2/BSA^{3/2})}$$  

where $RL_m = R/\mu BSA^{-3/2}$ is the viscosity-scaled resistance index.

4.5. Non-dimensional numbers in the clinical setting

Flow pulsatility has been shown to be beneficial for organ perfusion [27] as well as healthy vascular gene expression and remodelling; however, the energetic cost of delivering a constant CO is higher in pulsatile flow than in non-pulsatile flow [26,72]. In the healthy circulatory system, the pulsatile component of ventricular load operates optimally and constitutes a small fraction of the total power requirement of the ventricle (approx. 10%). However, many diseases, including congestive heart failure, ventricular hypertrophies, adverse vascular remodelling and congenital heart disease, as well as cardiovascular device performance (neonatal CPB, ventricle assist devices), are closely associated with increased pulsatile workload on the ventricle. Changes in the cardiovascular system that accompany such diseases were displayed as parametric variations in figure 3.

From the clinical perspective, surface-area scaling is commonly adopted for human physiology [35]; therefore, the PLI can be established as a size-invariant measure of cardiovascular performance. If quarter-power scaling is considered for humans, then a 10-fold variation in size (e.g. 10–100 kg) introduces a 30% variation in PLI between the two extremes of human size, which is still smaller than the variations in PLI introduced by disease conditions [18]. Therefore, the $PLI \propto M^{-1/6}$ relation that occurs in quarter-power scaling would only cause a weak body-size dependence of PLI when used as a performance index.

As discussed above, decay and propagation numbers correspond directly to the physical properties of the arterial system. Arterial pressure decay time is commonly associated with WK properties, which are determined by total arterial compliance and peripheral resistance. Characteristic arterial decay time ($\tau = RC$) is calculated by estimating the negative of the slope of the diastolic pressure decay curve [73]. We
note that $\delta$ is different from the $RC$ time used in the estimation of total arterial compliance [73]. The value of $RC$ time varies with the size of the subject, whereas $\delta$ is a scale-invariant index that indicates the WK performance, as it reflects the synchronization of heart rate with the decay properties of the vascular system. Similarly, $\psi$ is shown to be proportional to the normalized transit time of pulse waves from the aortic inlet to the microcirculation boundary, where transit time was estimated using the foot-to-foot method [74].

Figure 7a,b shows that there is a strong one-to-one association between the normalized decay time and normalized transit time, estimated from the simulated pressure waves and the non-dimensional $\delta$ and $\psi$ numbers given as inputs to the system. Estimated measures can be used directly to determine the state of the cardiovascular system, the ventricular–arterial coupling and the arterial pressure waves by referring to figures 3–5, without having to determine individual values of compliance and inertance.

4.6. Further limitations

The numerical model presented in this paper is the minimum model that is able to capture the transmission and WK properties of the circulation system. In the future, a distributed model of the circulation system may be tested against the present results. Linear models sufficiently capture the dynamics of circulation, however several advanced features may be included in future investigations. For example, nonlinear elasticity and viscoelasticity of the vessels, and inertial power losses, were not incorporated in the present model. Possible impacts of nonlinear effects are discussed in §4.2. Another limitation was that CO and flow waveforms were fixed. For advanced analysis of arterio-ventricular coupling, ventricular time-varying elastance or multi-physics models could be added [75–77].

At present, our model focused on a one inlet–one outlet transmission-WK model. However, vascular systems feature more complex branches with multiple inlets and outlets (e.g. carotid bifurcation (one inlet, two outlets), aortic arch (one inlet, four outlets), venous confluence (two inlets, one outlet) and hepatic confluence (three inlets, one outlet)), with fixed flow and pressure BCs as well as more complex BCs, which are determined by upstream and downstream vascular structures. For a more general power loss estimation, improved impedance calculations using the reduced-order, variable morphometric and fractal approaches [78–80] will be considered.

5. Conclusion

A complete set of non-dimensional numbers that govern both the steady and the pulsatile physics of cardiovascular systems are formulated. Non-dimensionalization of both lumped and distributed models of circulation is studied, covering the individual cardiovascular components as well as entire networks. Non-dimensional parameters obtained from a lumped arterial circulation model yield invariant characteristics, which were derived from common design principles of the circulatory systems of animals. We used the non-dimensionalization framework to elucidate the determinants of pulsatile load in the arterial system with a reduced number of variables. For the first time in the published literature, the complete optimal state maps of arterial circulation systems are calculated and the relative importance of pulsatile versus steady cardiovascular energy is quantified. The determinant non-dimensional variables governing decay and propagation characteristics can be acquired from the subject-specific routine clinical measurements, and would allow comparative assessment of optimal cardiovascular states. Extension of the presented framework towards the scaling of complex circulatory networks was also considered. Similitude of cardiovascular function across species will be useful in network optimization, as the only variable among the circuit candidates will be their ‘network topology’, allowing unbiased comparison of non-dimensional energetics and efficiency parameters.

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